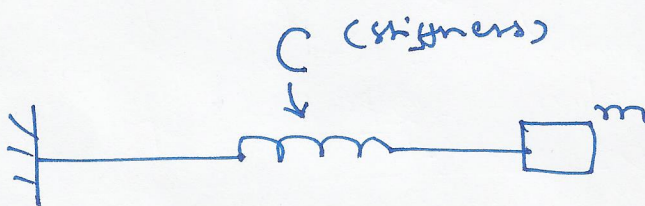


Classical realizations of Harmonic oscillator.

- Mass on a spring in the limit of small amplitude
- LC circuit for "small enough currents" (s.t., ckt. elem. are linear)
- Simple pendulum for small angle of oscillations.

Important properties of Harmonic oscillator

- f is independent of amplitude (in linear regime).
- Effects of several driving forces can be superimposed linearly.



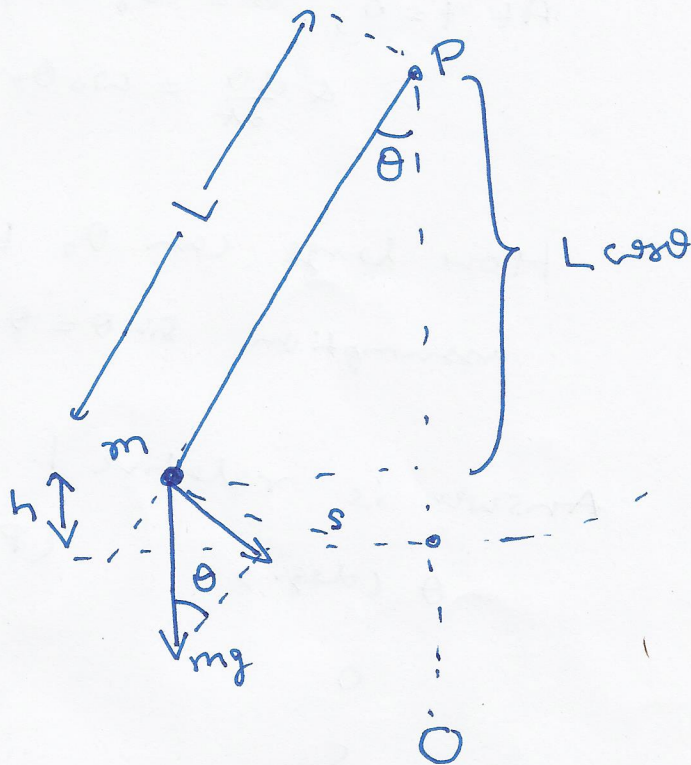
$$\vec{F} = -Cx\hat{x}$$

x is positive for stretch & negative for compression.

i.e., $M \frac{d^2x}{dt^2} = -Cx$.

$$\therefore x = A \sin(\omega_0 t + \phi)$$

$$\omega_0 = \left(\frac{C}{M}\right)^{1/2}; \quad \text{At } t=0, \quad x=x_0 = A \sin \phi; \quad \frac{dx}{dt} = v_0 = \omega_0 A \cos \phi.$$



$$s = L\theta$$

$$\therefore v = \frac{ds}{dt} = L \frac{d\theta}{dt} = L\dot{\theta}$$

$$a = \frac{d^2s}{dt^2} = L\ddot{\theta}$$

$\vec{F} = m\vec{a}$ reduces to

$$mg \sin\theta = -mL \frac{d^2\theta}{dt^2}$$

$$\sin\theta = \theta - \frac{\theta^3}{3!} + \frac{\theta^5}{5!} \dots$$

For small θ , $\sin\theta \approx \theta$.

$$\therefore \frac{d^2\theta}{dt^2} = -\frac{g}{L}\theta$$

$$\therefore \theta = \theta_0 \sin(\omega_0 t + \phi)$$

$$\text{where, } \omega_0 = \left(\frac{g}{L}\right)^{1/2}$$

$$f_0 = \frac{1}{2\pi} \sqrt{\frac{g}{L}}$$

Maximum value of amplitude is θ_0 .

$$\text{At } t=0, \theta = \theta_0 \sin \phi.$$

$$\& \frac{d\theta}{dt} = \omega_0 \theta_0 \sin \phi.$$

How large can θ_0 be so that assumption $\sin \theta = \theta$ is valid?

Answer is relative!

(Period $/ 2\pi\sqrt{L/g}$).

θ (deg.)

0	1.0000
5	1.0005
10	1.0019
15	1.0043
20	1.0077
30	1.0174
45	1.0396
60	1.0719

Can we deal with large amplitudes analytically?

Anharmonic oscillator

Let $\sin \theta = \theta - \frac{\theta^3}{6}$. (restrict to $\mathcal{O}(\theta^3)$.)

$$\therefore \frac{d^2 \theta}{dt^2} + \omega_0^2 \theta - \frac{\omega_0^2}{6} \theta^3 = 0.$$

$$\omega_0^2 = \frac{g}{L}.$$

Assume an approximate solⁿ.

$$\theta = \theta_0 \sin \omega t + \epsilon \theta_0 \sin 3\omega t.$$

ϵ is dimensionless constant
 $\epsilon \ll 1$ for $\theta_0 \ll 1$.

Why $3\omega t$?

$$\sin^3 x = \frac{3}{4} \sin x - \frac{1}{4} \sin 3x.$$

Thus θ^3 term in the diff. eqn.
 will generate from the cube of $\sin \omega t$
 a term in $\sin 3\omega t$.

$$\omega, \epsilon = ?$$

Note: $\epsilon \sin 3\omega t$ in trial solⁿ.
 will generate a term $\sim \sin 9\omega t$!

$$\ddot{\theta} = -\omega^2 \theta_0 \sin \omega t - 9\omega^2 \epsilon \theta_0 \sin 3\omega t.$$

$$\theta^3 = \theta_0^3 (\sin^3 \omega t + 3\epsilon \sin^2 \omega t \sin 3\omega t + \dots)$$

Use $\sin^3 x \equiv \frac{3}{4} \sin x - \frac{1}{4} \sin 3x.$

discarding terms $\mathcal{O}(\epsilon^2)$ & $\mathcal{O}(\epsilon^3).$

$$\ddot{\theta} = -\omega^2 \theta_0 \sin \omega t - 9\omega^2 \epsilon \theta_0 \sin 3\omega t$$

$$\omega_0^2 \theta = +\omega_0^2 \theta_0 \sin \omega t + \omega_0^2 \epsilon \sin 3\omega t.$$

$$-\frac{1}{6} \omega_0^2 \theta^3 = -\frac{3\omega_0^2 \theta_0^3}{24} \sin \omega t + \frac{\omega_0^2 \theta_0^3}{24} \sin 3\omega t.$$

$$-\frac{\omega_0^2 \theta_0^3}{2} \epsilon \sin^2 \omega t \sin 3\omega t.$$

coeff. of $\sin \omega t$ term must add up to zero,

$$\Rightarrow -\omega^2 + \omega_0^2 - \frac{3}{24} \omega_0^2 \theta_0^2 = 0.$$

$$\Rightarrow \omega^2 = \omega_0^2 \left(1 - \frac{1}{8} \theta_0^2\right).$$

$$\therefore \omega = \omega_0 \sqrt{1 - \frac{1}{8} \theta_0^2}.$$

$$\approx \omega_0 \left(1 - \frac{1}{16} \theta_0^2\right).$$

e.g., for $\theta_0 = 0.3 \text{ rad.}$,

$$\frac{\Delta \omega}{\omega} \approx -10^{-2}.$$

$$\Delta \omega = \omega - \omega_0.$$

\therefore Freq. of pendulum depends on amplitude for large amplitudes.

coeff. of $\sin 3\omega t$ term upto $\mathcal{O}(\theta_0^2)$,

$$-9\omega^2\epsilon + \omega_0^2\epsilon + \frac{\omega_0^2}{24}\theta_0^2 = 0.$$

upto $\mathcal{O}(\theta_0^2)$,

$$\omega^2 \approx \omega_0^2.$$

$$\therefore \epsilon \approx \frac{\theta_0^2}{192}$$

For $\theta_0 = 0.3 \text{ rad.}$,
we have, $\epsilon \approx 10^{-3}$.

Analogy

Harmonic oscillator & LC circuit

Voltage across the capacitance C is,

$V_C = \frac{Q}{C}$. (where Q is the charge on the capacitance)

Current in the ckt. in series with capacitance is,

$$I = -\frac{dQ}{dt} \quad \text{or, } Q = -\int I dt.$$

The voltage across the inductor L is,

$$V_L = -L \frac{dI}{dt}.$$

In LC ckt.,

$$-L \frac{dI}{dt} + \frac{Q}{C} = 0 = L \frac{d^2Q}{dt^2} + \frac{Q}{C}.$$

$$Q \leftrightarrow x$$

$$L \leftrightarrow M$$

$$\frac{1}{C} \leftrightarrow C_{\text{spring}}.$$

$$Q = Q_0 \sin(\omega_0 t + \phi)$$

$$\omega_0 = \left(\frac{1}{LC}\right)^{1/2}.$$